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ON A LIPSCHITZ STABILITY PROBLEM FOR p-LAPLACIAN BESSEL EQUATION

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ABSTRACT. In this study, we are enunciative of some asymptotic expansions and reconstruction formulas for inverse nodal problem of p-Laplacian Bessel equation. Furthermore, Lipschitz stability problem for this equation with Dirichlet boundary conditions is solved. And, it is also proved that the space of all potential functions w is homeomorphic to the partition set of all asymptotically equivalent nodal sequences induced by an equivalence relation.

1. INTRODUCTION

Let us consider the below p-Laplacian Bessel eigenvalue problem;

$$-\left(u^{'(p-1)}\right)' = (p-1)\left(\lambda - w_0(x) - \frac{l(l+1)}{x^2}\right)u^{(p-1)}, \ 1 \le x \le a,$$
(1.1)

with Dirichlet conditions

$$u(1) = u(a) = 0, (1.2)$$

where l = 0, 1, 2, ..., a, p > 1 are constants, λ is a spectral parameter; $w_0(x) \in L^2[1, a]$ is a real valued function and $u^{(p-1)} = |u|^{(p-1)} sgnu$ (see [1], [2]). Normally, the equation (1.1) is considered by a condition at the origin. In this instance, the problem becomes singular and it is not easy to solve inverse nodal problem for these type equations in p-Laplacian case. Therefore, we will study the Lipschitz stability of inverse nodal problem for p-Laplacian Bessel equation on a smooth interval.

Uniqueness and reconstruction problems of p-Laplacian Bessel equation have been considered in some works (see [2], [3]) just left a stability issue is worth considering and undone for eigenvalue problem (1.1)-(1.2). In an exact solution of inverse problems, the questions of existence, uniqueness, stability and construction are to be taken into account. The query of existence and uniqueness is of major

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importance in testing the assumption behind any mathematical model. If the response in the uniqueness question is no, then we know that even perfect data do not contain enough data to recover the physical quantity to be estimated. In the query of stability, we should decide whether the solution depends continuously on the data. Stability is essential if we want to be sure that a variation of the given data in a sufficiently small range leads to an arbitrarily small modification in the solution. This notion was firstly defined by Hadamard in 1902 (see [4]). Due to this significant reason, we want to deal with a stability issue for the problem (1.1)-(1.2).

As clearly seen, equation (1.1) reduces to the following equation

$$-u'' + \left(w_0 + \frac{l(l+1)}{x^2}\right)u = \lambda u,$$
(1.3)

for p = 2 and it is known as Bessel equation which is obtained by separation of variables in the 3D radial Schrödinger equation. In case of a wave function with spherical symmetry, the wave equation can be separated using spherical coordinates, and the equation for the radial component becomes (1.3) where λ is a constant refers to eigenvalue of the problem, physically proportional to the energy of the particle under consideration, w_0 is proportional to the potential energy and l is a positive integer or zero. Throughout this study, we will use this equality $w_0 + \frac{l(l+1)}{x^2} = w$, briefly. Inverse problems in spectral theory are divided into two parts; inverse eigenvalue problem and inverse nodal problem. Inverse eigenvalue problem is constructing operator by using some spectral datas as spectrum, norming constants. The first study is given by Ambarzumyan in 1929 about inverse eigenvalue problems [5]. Later, many researchers obtained some important results for different type operators. Inverse eigenvalue problem for classical Bessel equation was studied by various authors (see [6], [7], [8], [9], [10], [11], [12], [13], [14], [15]).

In 1988, McLaughlin [16] proposed a new way to construct the Sturm-Liouville operator. This effective technique is called inverse nodal problem. These problems consist in recovering operators from given zeros of their eigenfunctions (nodal points). Physically, this is related to find, e.g., the density of a string or a beam from the zero-amplitude positions of their eigenvibrations. She is admitted to be the first mathematician to consider this kind of inverse problem. Independently, Shen, studied the relation between density function of string equation and nodal points [17]. Afterwards, inverse nodal problem has been kept in view by several authors (see [18], [19], [20]).

Assume that $X_n = \{x_j^n\}_{j=1}^{n-1}$ are the zeros of the eigenfunction $u_n(x)$ related to the eigenvalues $\{\lambda_n\}$ of the problem (1.1)-(1.2). Nodal length of this problem is denoted by $l_j^n = x_{j+1}^n - x_j^n$ for j = 1, 2, ..., n-1. Using these type nodal datas, some uniqueness and reconstruction results for different type of operators have been expressed by several authors (see [21], [22], [23], [24]).

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Equation (1.1) gives us the following p-Laplacian Sturm-Liouville eigenvalue problem in case of w(x) = 0,

$$-\left(u^{'(p-1)}\right)' = (p-1)\lambda u^{(p-1)}, u(1) = u(a) = 0.$$
(1.4)

where the eigenvalues of the problem (1.4) associated eigenfunctions $u_n(x)$ are countably infinite real and simple [25]. Inverse and stability problems for (1.4) were studied many times in literature (see [26], [27], [28], [29], [30], [31]). Besides, inverse problems for different types of p-Laplacian operators have been proved by various mathematicians (see [32], [33], [34], [35], [36]).

To say something about the stability of the inverse nodal problem for the problem (1.1)-(1.2), we should define the function S_p that is the solution of the initial value problem (see [25], [32], [33])

$$-\left(S_{p}^{\prime(p-1)}\right)' = (p-1)S_{p}^{(p-1)}, S_{p}(0) = 0, \ S_{p}'(0) = 1.$$
(1.5)

 S_p and S'_p are periodic functions which provide the identity

$$|S_p(x)|^p + |S'_p(x)|^p = 1$$

for any $x \in \mathbb{R}$ which is similar with the known trigonometric identity $\sin^2 x + \cos^2 x = 1$ where $S_p(x)$ is called generalized sine function. These generalized functions are p-analogues of classical sine and cosine functions. It is well known that

$$\pi_p = 2 \int_0^1 \frac{dt}{(1 - t^p)^{\frac{1}{p}}} = \frac{2\pi}{p \sin\left(\frac{\pi}{p}\right)}$$

is the first zero of S_p (see [33]). Presently, we will give some important properties of S_p .

Lemma 1.1. [25], [33] a) For $S'_p \neq 0$,

$$\left(S'_p\right)' = -\left|\frac{S_p}{S'_p}\right|^{p-2} S_p$$

b)

$$\left(S_p S_p^{\prime (p-1)}\right)' = \left|S_p'\right|^p - (p-1)S_p^p = 1 - p\left|S_p\right|^p = (1-p) + p\left|S_p'\right|^p.$$

This study is organized as follows. In section 2, we express some asymptotic formulas for eigenvalues, nodal parameters and potential function of the problem (1.1)-(1.2) by using modified Prüfer substitution [2]. In section 3, we define a metric to clarify Lipschitz stability of inverse nodal problem for p-Laplacian Bessel equation. Finally, we give some conclusions in section 4.

2. Asymptotic formulas for p-Laplacian Bessel equation

In this section, we recur some features of p-Laplacian Bessel equation (1.1) with the conditions (1.2) which were solved in [2]. In accordance with this purpose, we need to define the following modified Prüfer substitution as

$$u(x) = c(x)S_p\left(\lambda^{1/p} \theta(x,\lambda)\right), \qquad (2.1)$$
$$u'(x) = (l+1)\lambda^{1/p}c(x)S'_p\left(\lambda^{1/p} \theta(x,\lambda)\right),$$

or

$$\frac{u'(x)}{u(x)} = (l+1)\lambda^{1/p} \frac{S'_p\left(\lambda^{1/p} \ \theta(x,\lambda)\right)}{S_p\left(\lambda^{1/p} \ \theta(x,\lambda)\right)},\tag{2.2}$$

where c(x) is amplitude and $\theta(x)$ is Prüfer variable [37]. Differentiating each side of the above equation (2.2) with respect to x and considering Lemma 1.1, it can be obtained as [2]

$$\theta'(x) = l + 1 + \left[-(l+1) + (l+1)^{1-p} - \frac{(l+1)^{1-p}}{\lambda} \left(w_0(x) + \frac{l(l+1)}{x^2} \right) \right] S_p^p \left(\lambda^{1/p} \ \theta(x,\lambda) \right)$$
(2.3)

This equality plays an important role throughout this study. Now, we are ready to express asymptotic estimations of nodal parameters and potential function for the problem (1.1), (1.2).

Theorem 2.1. [2] The eigenvalues of the problem (1.1)-(1.2) have the form

$$\begin{split} \lambda_n^{1/p} &= \frac{n\pi_p}{\tilde{l}\,(a-1)} + \frac{(l+1)^{1-p}\tilde{l}^{p-2}(a-1)^{p-2}}{p(n\pi_p)^{p-1}} \int_1^a \left\{ w_0(s) + \frac{l(l+1)}{s^2} \right\} ds + O\left(\frac{1}{n^{p-1}}\right), \\ as \ n \to \infty \ where \ \tilde{l} &= (l+1)\left(1 - \frac{1}{p} + \frac{1}{p(l+1)^p}\right). \end{split}$$

Theorem 2.2. [2] The nodal parameters of the problem (1.1)-(1.2) satisfy following equalities

$$\begin{split} x_{j}^{n} &= 1 + \frac{j(a-1)\tilde{l}}{(l+1)n} - \frac{j\tilde{l}^{p}(a-1)^{p}}{p(l+1)^{p}n^{p+1}\pi_{p}^{p}} \int_{1}^{a} w(s)ds \\ &+ \int_{1}^{x_{j}^{n}} S_{p}^{p}ds - \frac{1}{(l+1)^{p}} \int_{1}^{x_{j}^{n}} \left\{ 1 - \frac{\tilde{l}^{p}(a-1)^{p}w(s)}{(n\pi_{p})^{p}} \right\} S_{p}^{p}ds + O\left(\frac{j}{n^{p+1}}\right), \end{split}$$

and

$$l_j^n = \frac{a-1}{n} + \frac{(a-1)^{p\tilde{l}^{p-1}}}{p\tilde{l} \ (l+1)^{p-1} \ (n\pi_p)^p} \int_{x_j^n}^{x_{j+1}^+} w(t)dt + O\left(\frac{1}{n^p}\right),$$

respectively, as $n \to \infty$.

Theorem 2.3. [2] Let $w \in L^2[1, a]$. Then

$$w(x) = \lim_{n \to \infty} p(l+1)^{p-1} \lambda_n \left(\frac{\tilde{l} \lambda_n^{\frac{1}{p}}}{\pi_p} l_j^n - 1 \right),$$

where $x \in (1, a)$ and $j = j_n(x) = \max\{j : x_j^n < x\}.$

3. MAIN RESULTS ABOUT LIPSCHITZ STABILITY

Here, we want to clarify Lipschitz stability problem for *p*-Laplacian Bessel equation. Lipschitz stability is about a continuity between two metric spaces. To denote this continuity, we will handle a homeomorphism between these metric spaces. Stability problems for different operators were studied by many authors (see [38], [39], [40], [41], [42]). The procedure that we have applied in the proof of Lipschitz stability is similar to the study [38].

Let define Ω_{dif} and Σ_{dif} by

$$\begin{aligned} \Omega_{dif} &= \left\{ w \in L^2 \left[1, a \right] \right\}, \\ \Sigma_{dif} &= \left\{ X = \left\{ x_k^n \right\} : X \text{ is the nodal set associated with some } w \in \Omega_{dif} \right\}. \end{aligned}$$

We will denote that Ω_{dif} and Σ_{dif} are homeomorphic to each other. Therefore, when \overline{X} is the nodal set associated with \overline{w} and \overline{X} is close to X in Σ_{dif} , then \overline{w} is close to w in Ω_{dif} . Hence, inverse nodal problem for p-Laplacian Bessel equation is Lipschitz stable. Herein, we use $L^m(1, a)$ $(m \ge 1)$ for Ω_{dif} . Let

$$S_n^m(X,\overline{X}) = \frac{(l+1)^{p-1}\pi_p^p}{\tilde{l}^p(a-1)^{p+1-\frac{1}{m}}} n^{p+1-\frac{1}{m}} \left[\sum_{k=0}^{n-1} \left| l_k^n - \bar{l}_k^n \right|^m \right]^{\frac{1}{m}}, \quad (3.1)$$

where $l_k^n = x_{k+1}^n - x_k^n$, $\overline{l}_k^n = \overline{x}_{k+1}^n - \overline{x}_k^n$ and $m \ge 1$. Define the metric and a pseudometric on Σ_{dif}

$$d_0^m(X,\overline{X}) = \overline{\lim_{n \to \infty}} S_n^m(X,\overline{X}),$$

and

$$d_{\Sigma_{dif}}^m(X,\overline{X}) = \overline{\lim_{n \to \infty}} \frac{S_n^m(X,\overline{X})}{1 + S_n^m(X,\overline{X})},$$

respectively. If we set a relation as $X \sim_m \overline{X}$ iff $d^m_{\Sigma_{dif}}(X,\overline{X}) = 0$, then \sim_m is an equivalence relation on Σ_{dif} and $d^m_{\Sigma_{dif}}$ would be a metric for the partition set $\Sigma^*_{dif} = \Sigma_{dif} / \sim_m$.

Lemma 3.1. $d_{\Sigma_{dif}}^m(.,.)$ is a pseudometric on Σ_{dif} .

Proof: The proof can be expressed easily by considering the similar way with in [25].

Lemma 3.2. Let $X, \overline{X} \in \Sigma_{dif}$. Then,

(a): The interval I_{n,k} between the points xⁿ_k and xⁿ_k has length O(n^{-p}).
(b): The inequality |j_n(x) - j
_n(x)| ≤ 1 holds when n is sufficiently large for all x ∈ (1, a).

Proof:

where A =

(a) Using the similar way with in [38] and the asymptotic expansions of nodal points yield

$$\begin{aligned} |I_{n,k}| &= |x_k^n - \overline{x}_k^n| \\ &\leq |x_k^n - A| + |A - \overline{x}_k^n| \\ &= O(n^{-p}) + O(n^{-p}) \\ &= O(n^{-p}), \\ \frac{l+1}{\tilde{l}} + \frac{k(a-1)}{n} - \frac{l+1}{p\tilde{l}} + \frac{1}{p\tilde{l}(l+1)^{p-1}}. \end{aligned}$$

(b): It can be proved by using analogous method with [38].

Theorem 3.1. $d_{\Sigma_{dif}}^m$ is a metric on the space Σ_{dif} / \sim_m for any of $m \ge 1$. Intercalarly, the metric spaces $(\Omega_{dif}, \|.\|_m)$ and $(\Sigma_{dif} / \sim_m, d_{\Sigma_{dif}}^m)$ are homeomorphic to each other where \sim_m is an equivalence relation induced by $d_{\Sigma_{dif}}^m$.

Proof: It is sufficient to show that

$$\|w - \overline{w}\|_m = pd_0^m(X, \overline{X}).$$

By Theorem 2.3, we get

$$w(x) - \overline{w}(x) = \lim_{n \to \infty} \frac{p(l+1)^{p-1} (n\pi_p)^p}{\tilde{l}^p (a-1)^p} \left[\frac{n}{a-1} \left(l_{j_n(x)}^n - \overline{l}_{\overline{j}_n(x)}^n \right) \right],$$

for each $x \in (1, a)$. Therefore, we obtain

$$\begin{split} \|w - \overline{w}\|_{m} &\leq \frac{p(l+1)^{p-1}n(n\pi_{p})^{p}}{\tilde{l}^{p}(a-1)^{p+1}} \lim_{n \to \infty} \left\| l_{j_{n}(x)}^{n} - \overline{l}_{\overline{j}_{n}(x)}^{n} \right\|_{m} \\ &\leq \frac{p(l+1)^{p-1}\pi_{p}^{p}}{\tilde{l}^{p}(a-1)^{p+1}} \lim_{n \to \infty} \left[n^{p+1} \left\| l_{j_{n}(x)}^{n} - \overline{l}_{j_{n}(x)}^{n} \right\|_{m} \right].$$
(3.2)

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by Fatou's Lemma and the definition of norm on L_m . Here, Lemma 3.2. and Theorem 2.2. yield

$$n^{p+1} \left\| \bar{l}_{j_{n}(x)}^{n} - \bar{l}_{\bar{j}_{n}(x)}^{n} \right\|_{m} = n^{p+1} \left[\int_{1}^{a} \left| \bar{l}_{j_{n}(x)}^{n} - \bar{l}_{\bar{j}_{n}(x)}^{n} \right|^{m} dx \right]^{\frac{1}{m}}$$
$$= n^{p+1} \left[\sum_{k=0}^{n-1} \left| \bar{l}_{k+1}^{n} - \bar{l}_{k}^{n} \right|^{m} I_{n,k} \right]^{\frac{1}{m}}$$
$$= O\left(n^{1+\frac{1-p}{m}} \right) = o(1), \qquad (3.3)$$

and

$$n^{p+1} \left\| l_{j_n(x)}^n - \bar{l}_{j_n(x)}^n \right\|_m = n^{p+1} \left[\int_{1}^{a} \left| l_{j_n(x)}^n - \bar{l}_{j_n(x)}^n \right|^m dx \right]^{\frac{1}{m}}$$
$$= n^{p+1} \left[\sum_{k=0}^{n-1} \left| l_k^n - \bar{l}_k^n \right|^m l_k^n \right]^{\frac{1}{m}}$$
$$= n^{p+1-\frac{1}{m}} (a-1)^{\frac{1}{m}} \left[\sum_{k=0}^{n-1} \left| l_k^n - \bar{l}_k^n \right|^m \right]^{\frac{1}{m}}. \quad (3.4)$$

Considering (3.3) and (3.4) together in (3.2), we obtain

$$\begin{split} \|w - \overline{w}\|_{m} &\leq \frac{p(l+1)^{p-1} \pi_{p}^{p}}{\tilde{l}^{p} (a-1)^{p+1-\frac{1}{m}} \overline{\lim_{n \to \infty}} n^{p+1-\frac{1}{m}} \left[\sum_{k=0}^{n-1} \left| l_{k}^{n} - \overline{l}_{k}^{n} \right|^{m} \right]^{\frac{1}{m}} \\ &= p d_{0}^{m} (X, \overline{X}). \end{split}$$

Contrarily, using the above derivations

$$\begin{split} \|w - \overline{w}\|_{m} + o(1) &= \frac{p(l+1)^{p-1} \pi_{p}^{p}}{\tilde{l}^{p}(a-1)^{p+1}} n^{p+1} \left\| l_{j_{n}(x)}^{n} - \bar{l}_{\overline{j}_{n}(x)}^{n} \right\|_{m} \\ &\geq \frac{p(l+1)^{p-1} \pi_{p}^{p}}{\tilde{l}^{p}(a-1)^{p+1}} n^{p+1} \left\| l_{j_{n}(x)}^{n} - \bar{l}_{\overline{j}_{n}(x)}^{n} \right\|_{m} - O\left(n^{1+\frac{1-p}{m}}\right) \\ &= \frac{p(l+1)^{p-1} \pi_{p}^{p}}{\tilde{l}^{p}(a-1)^{p+1}} n^{p+1} \left[\sum_{k=0}^{n-1} \left| l_{k}^{n} - \bar{l}_{k}^{n} \right|^{m} l_{k}^{n} \right]^{\frac{1}{m}} - O\left(n^{1+\frac{1-p}{m}}\right) \\ &= \frac{p(l+1)^{p-1} \pi_{p}^{p} n^{p+1-\frac{1}{m}}}{\tilde{l}^{p}(a-1)^{p+1-\frac{1}{m}}} \left[\sum_{k=0}^{n-1} \left| l_{k}^{n} - \bar{l}_{k}^{n} \right|^{m} \right]^{\frac{1}{m}} - O\left(n^{1+\frac{1-p}{m}}\right). \end{split}$$

Hereby, as n approaches to infinity,

$$\|w - \overline{w}\|_m \ge pd_0^m(X, \overline{X}).$$

So, the proof is complete.

4. Conclusion

In this study, we have underlined the significance of the stability issue for inverse nodal problem of p-Laplacian Bessel equation. Then, some asymptotic expansions for eigenvalues, nodal parameters and reconstruction formula for potential function of the problem (1.1)-(1.2) have been expressed. Moreover, we have solved the Lipschitz stability problem for (1.1)-(1.2).

Competing Interests

The authors declare to have no competing interests.

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